# Subgroup meeting-08/24 Introduction of thermal transport



### Outline

- Introduction
- Phonons
- Specific heat of solid
- Conclusion
- Q and A

### Introduction



• Electronic charge carriers and Phonons.

 $\kappa_{total} = \kappa_e + \kappa_L$ 

( $\kappa$  is thermal conductivity)

### Phonons

#### Quasiparticle

quantization of the modes of lattice vibrations of periodic ,elastic crystal structure of solid.

#### Lattice waves(lattice vibration)



Moreover..

# Specific heat of solids

- $c = Q/m\Delta T$
- For fixed volume:  $C_v = \left(\frac{\partial E}{\partial T}\right)_{V^{+1}}$
- Three models
  Dulong-Petit model
  Bose-Einstein model
  Debye model

#### Three models



..... Law of Dulong – Petit, based on Maxwell-Boltzmann distribution

Bose-Einstein model

Debye model

Q: why specific heat of solid behaves differently at low and high temp?



A:Quantum effect emerges at low temp, and this effect is disturbed by thermal energy at high temp.

### **Dulong and Petit**

#### Maxwell-Boltzman distribution function

E =  $k_BT$ .... ( $k_B$ : Boltzman const) three dimension, E =  $3k_BT$  $C_v = \partial E/\partial T = \partial Q/\partial T$  $C_v = 3k_B$ 

#### only valid at high temp regime!

#### specific heat of solid must be calculated in different ways at low and high temps respectivel

### **Bose-Einstein model**

#### Assumption

- Three dimension harmonic oscillator
- Vibrate independently
- same frequency







- Mistake 1: only one atomic vibration frequency  $(\omega_{\rm E})$
- Mistake 2: no interaction between individual phonons.

#### Assumption

- Similar to black body
- Vibrate dependently as elastic standing waves (differ from Bose-Einstein model)
- Frequency o~v



frequency (Einstein model) Frequency distribution (Debye model)

• atomic vibrations as phonons in a box

$$\lambda_{n} = \frac{L}{2n}$$
$$E_{n} = \hbar \omega_{n} = \frac{hv}{\lambda_{n}} = \frac{2nhv}{L}$$

Bose-Einstein statistics

$$\overline{N(E_n)} = \frac{1}{e^{E/kT} - 1}$$

• The energy in solid

$$U = \sum E_n \overline{N(E_n)}$$

$$U = \sum_{K} \sum_{P} \frac{\hbar \omega_{K,P}}{\exp\left(\frac{\hbar \omega}{k_{R}T}\right) - 1}$$

$$U = \sum_{p} \int \frac{\hbar \omega_{K}, p}{\exp\left(\frac{\hbar \omega}{k_{B}T}\right) - 1} D_{p}(\omega) d\omega \quad D(\omega) = \text{density of modes/states}$$

#### 1-D normal vibration mode



$$x_{s} = x(0) \exp(-i\omega_{K,p}t) \sin sKa$$
$$K = \frac{\pi}{L}, \frac{2\pi}{L}, \dots, \frac{(N-1)\pi}{L}$$

### frequency distribution

#### periodic boundary



- sin(sKa),cos(sKa)
- boundary conditions
  NKa = 2nπ

#### For K=0 , K=8π/8a , sin(sKa) is meaningless

periodic boundary condition



 $\mathbf{\Lambda}$ 

**K**-space

3-D normal vibration density of states

• The allowed K is 
$$\left(\frac{2\pi}{L}\right)^3$$

The number of allowed states

$$N = \left(\frac{4\pi K^{3}}{3}\right) / \left(\frac{2\pi}{L}\right)^{3}$$
$$D(\omega) = \frac{dN}{d\omega} = \left(\frac{VK^{2}}{2\pi^{2}}\right) \left(\frac{dK}{d\omega}\right)$$

The approximation of Debye model

$$D(\omega) = \frac{V\omega^2}{2\pi^2\nu^3} \qquad \omega = \nu K$$

Debye frequency

$$\omega_{\rm D}{}^3 = \frac{6\pi^2\nu^3N}{V}$$

• The energy in solid

$$U = 3 \times \int_{0}^{\omega_{D}} \left( \frac{V \omega^{2}}{2\pi^{2} \nu^{3}} \right) \left( \frac{\hbar \omega}{\exp\left(\frac{\hbar \omega}{k_{B}T}\right) - 1} \right) d\omega$$

• Debye temperature is defined as  $\theta$ 

$$\boldsymbol{\Theta} = \frac{\hbar\omega}{k_{\rm B}T} = \frac{\nu}{k_{\rm B}T} (\frac{6\pi^2 N}{V})^{1/3}$$

• The energy in solid (Debye model)  $U = 3Nk_{B}TD_{3}(\frac{\theta}{T}) \qquad C_{V} = 9Nk_{B}(\frac{\theta}{T})^{3}\int_{0}^{x_{D}}\frac{x^{4}e^{x}}{(e^{x}-1)^{2}}dx$ 



#### Lattice thermal transport: (wave-particle duality description)

Classical mechanics: wave-like phenomena Quantum mechanics: particle-like properties

# Thanks for your attention!!

# Q and A